

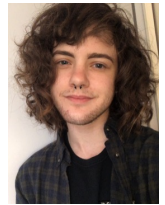
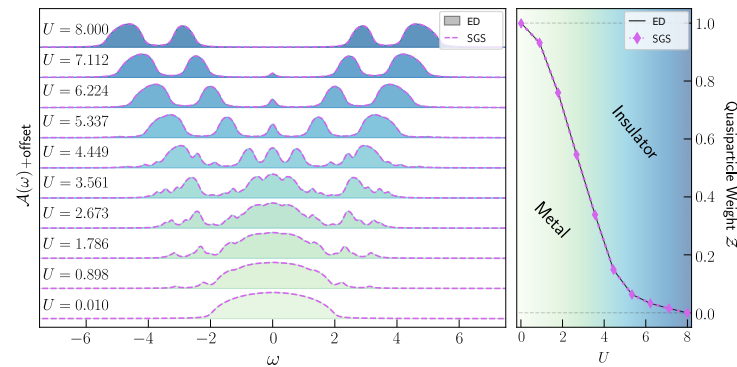
A quantum computing approach to efficiently simulating correlated materials using impurity models and DMFT

Alexander (Lex) Kemper



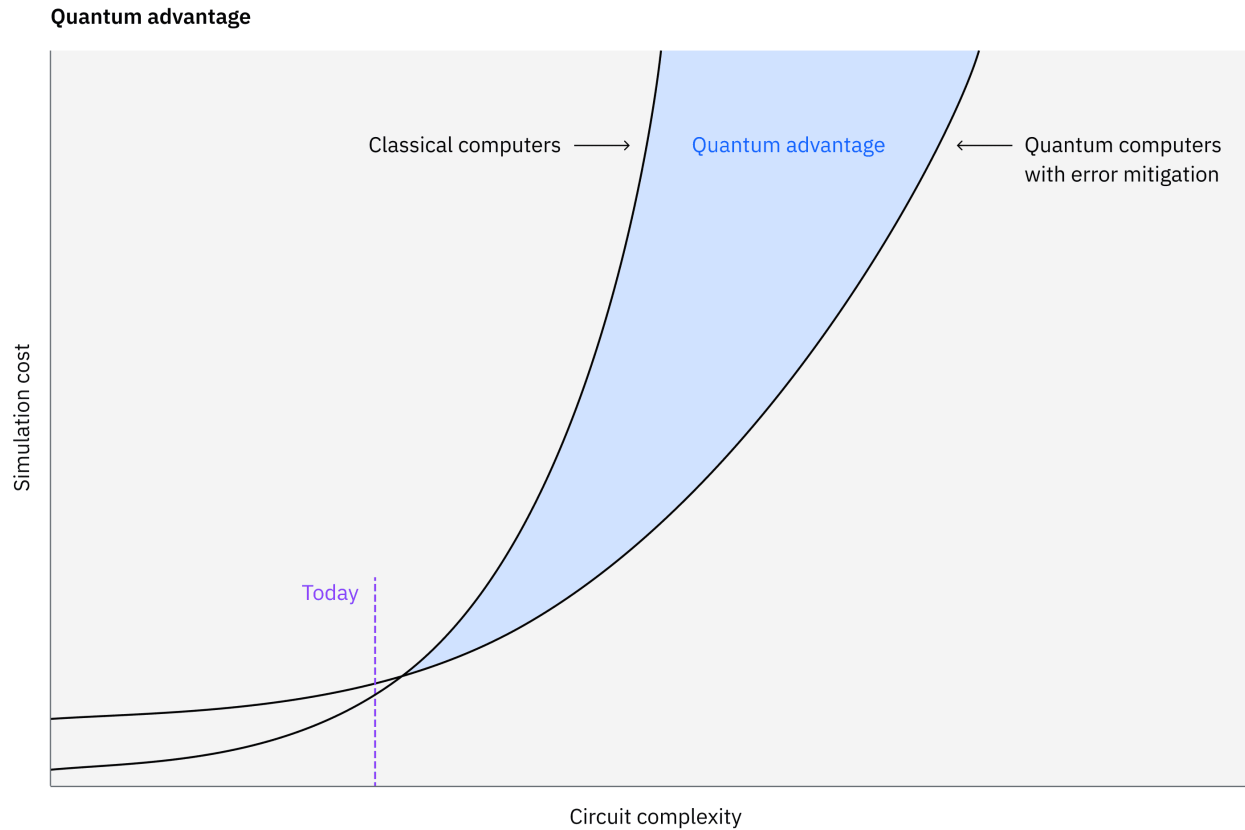
Department of Physics
North Carolina State University
<https://go.ncsu.edu/kemper-lab>

CCDS25
08-09-2025



[N. Hogan, E. Kökcü, T. Steckmann, L. Doak, C. Mejuto-Zaera, D. Camps, R. van Beeumen, W.A. de Jong, A.F. Kemper, 2508.05738](#)

Quantum computers in 2025



<https://www.ibm.com/quantum/blog/quantum-advantage-era>

Quantum computers in 2025

Bespoke quantum simulator



Digital algorithms

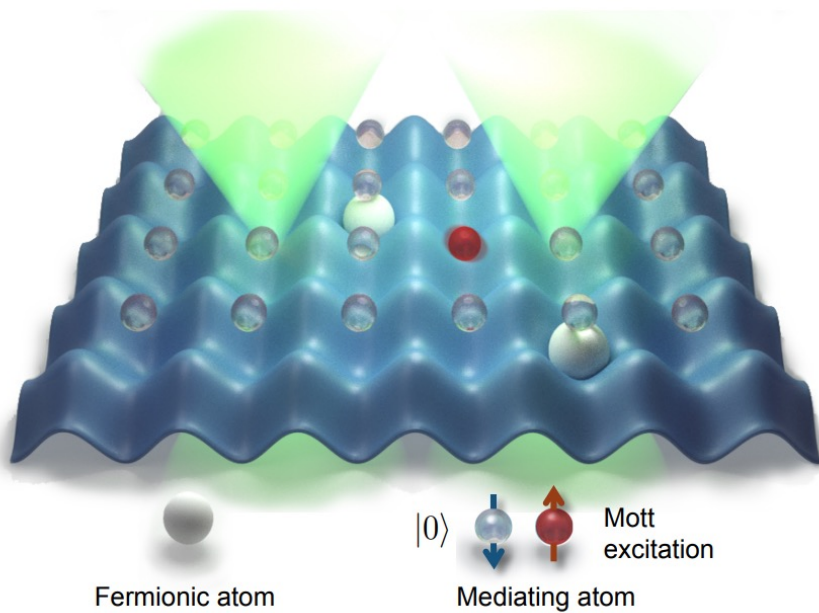
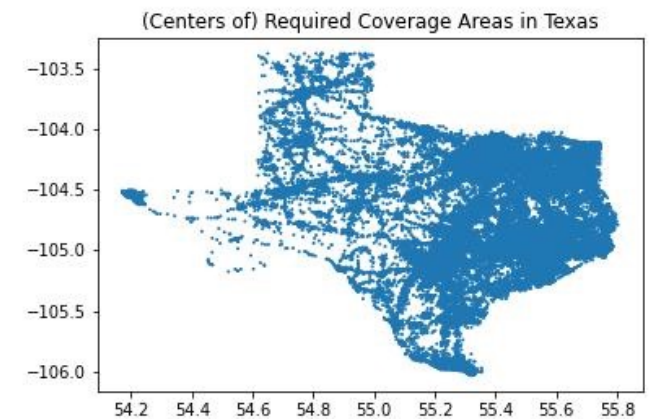
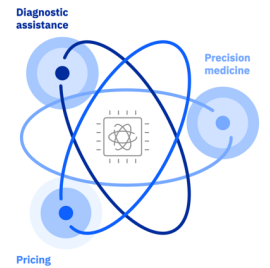


Figure 1
Quantum computers may enable three key healthcare use cases that reinforce each other in a virtuous cycle. For instance, accurate diagnoses enable precise treatments, as well as a better reflection of patient risks in pricing models.

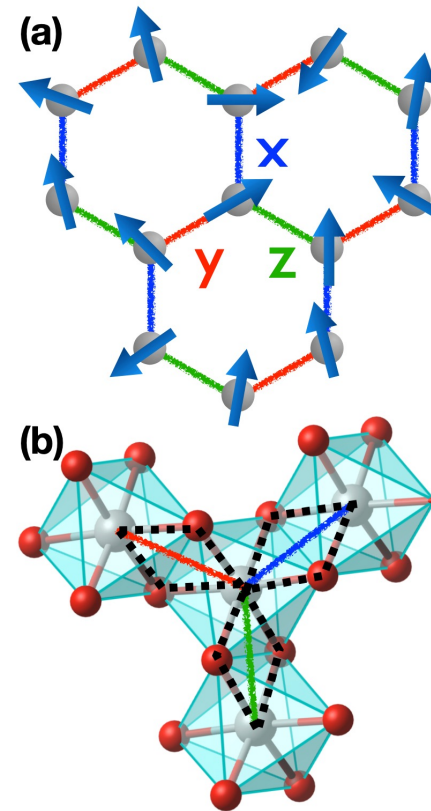
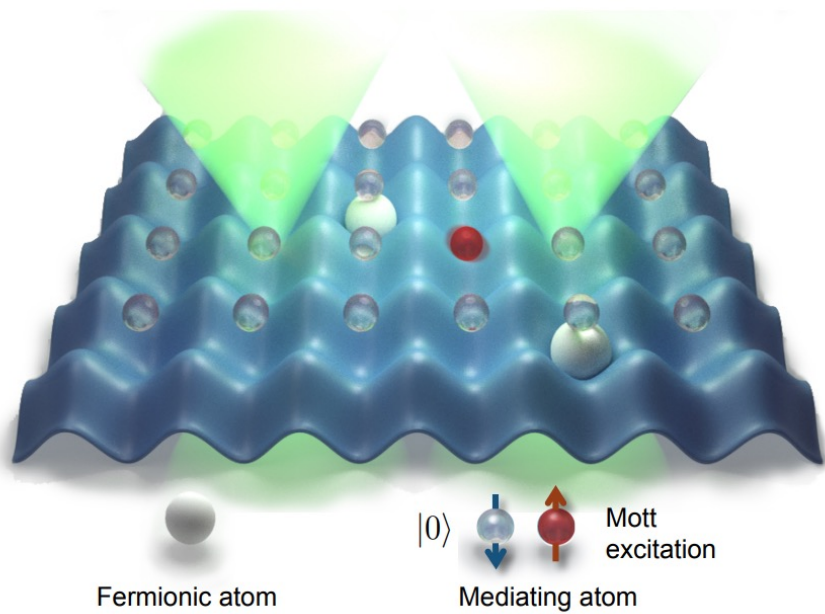


Quantum computers in 2025

Bespoke quantum simulator



Digital algorithms



Lex' definition of quantum advantage: when I can use a quantum computer to answer a question relevant to condensed matter physicists.

To get to a quantum advantage, we need a problem that is

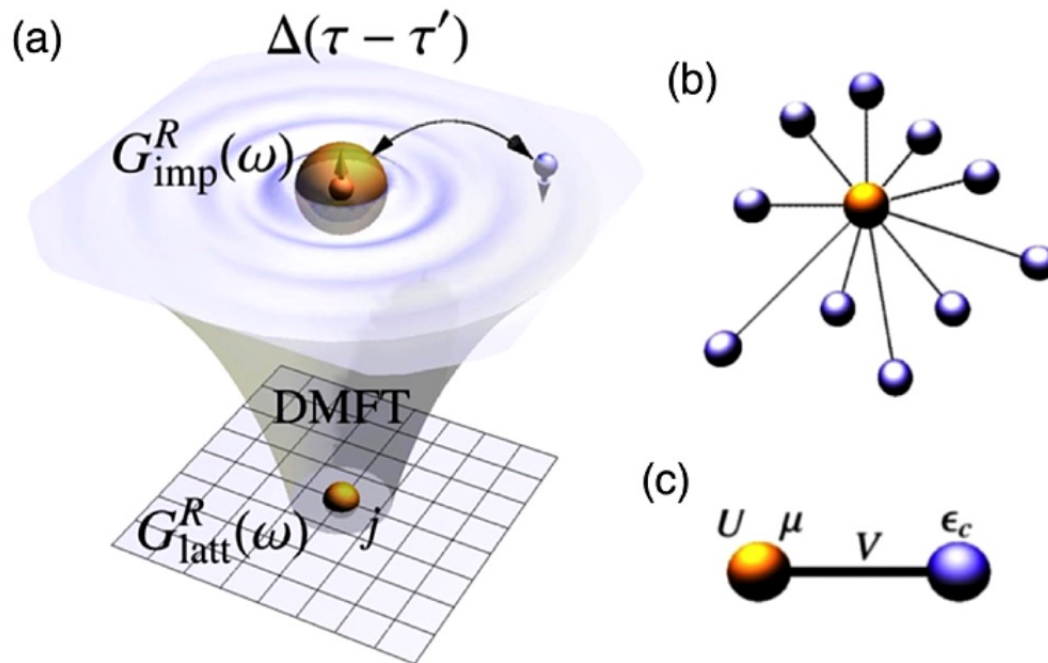
- Relevant/interesting
- Can be used to interface with non-QC folks
- Runs on a few qubits (< 100)
- Doesn't require long qubit coherence times

Quantum computers in 2025

Lex' de
comp
phys

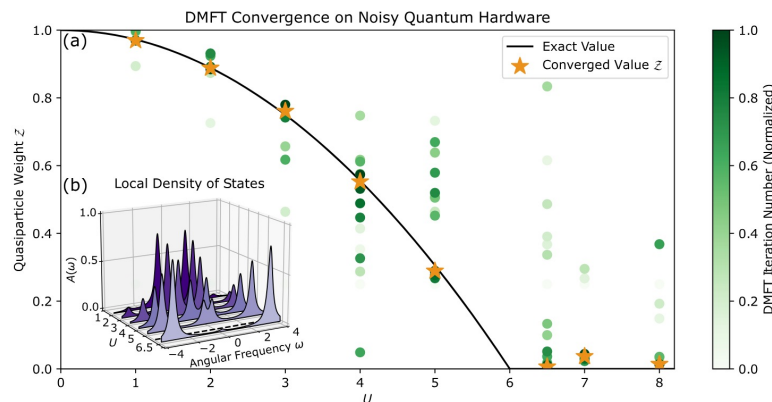
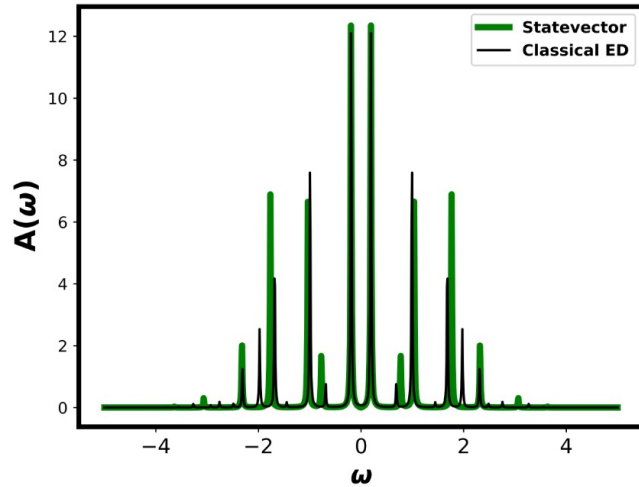
To

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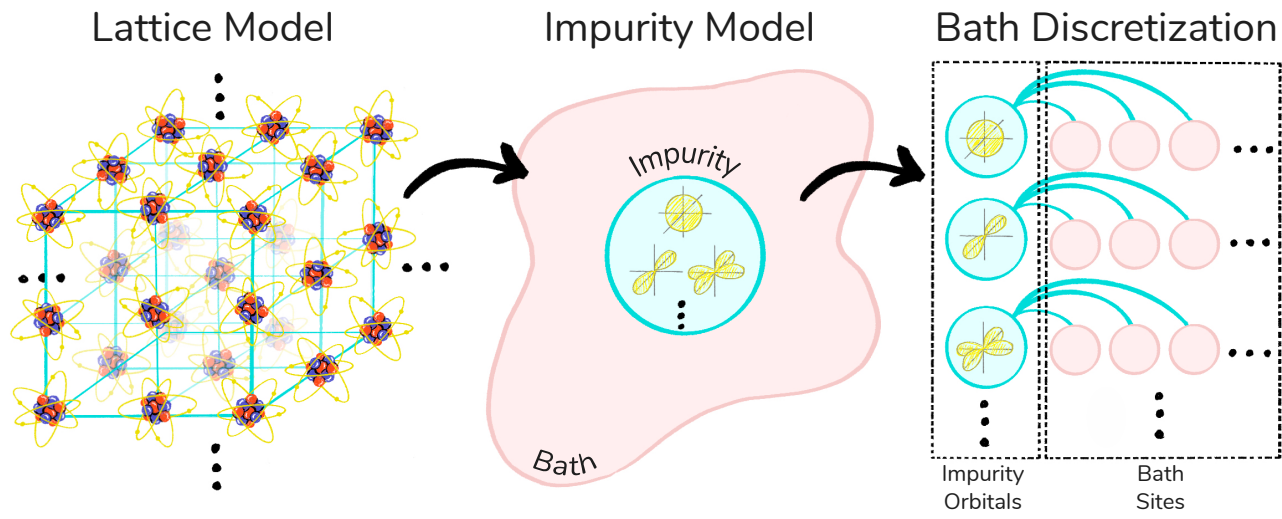
Kreula EPJ Quant. Tech. (2016)

DMFT on Quantum Computers



- Bauer PRX (2016)
- Kreula, EPJ QT (2016)
- Rungger arXiv (2019)
- Keen QST (2020)
- Besserve PRB (2022)
- Jamet arXiv (2022)
- Steckmann PRR (2023)
- Nie PRL (2024)
- Selisko arXiv (2024)
- Greene-Diniz Quantum (2024)
- Jamet APL Quantum (2025)

DMFT on Quantum Computers



$$\mathcal{G}_{ij}(t) = -i\theta(t) \langle \psi | \{c_i(t), c_j^\dagger\} | \psi \rangle$$

Correlation functions

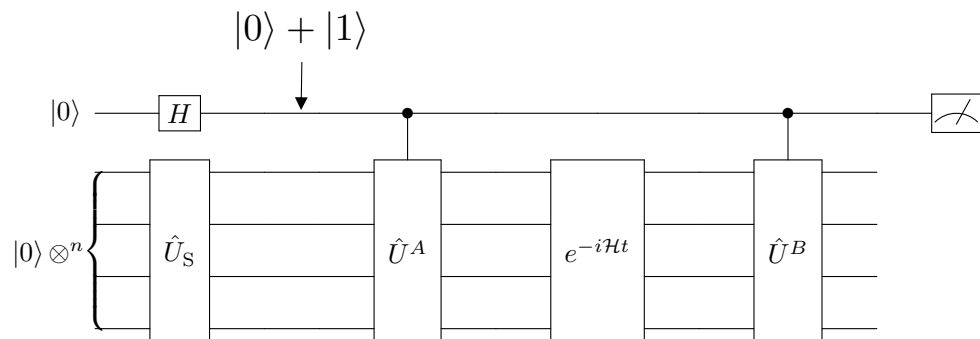
The Rules

1. You may only use unitary operations
2. More qubits = bad
3. More operations = bad
4. Complex qubit operations = bad
5. Your results will be very noisy anyway

$$e^{iE_0 t} \langle \phi_0 | B e^{-i\mathcal{H}t} A | \phi_0 \rangle$$

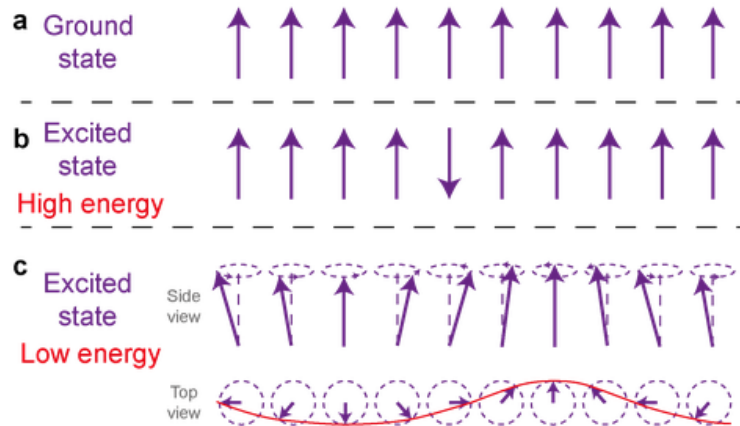
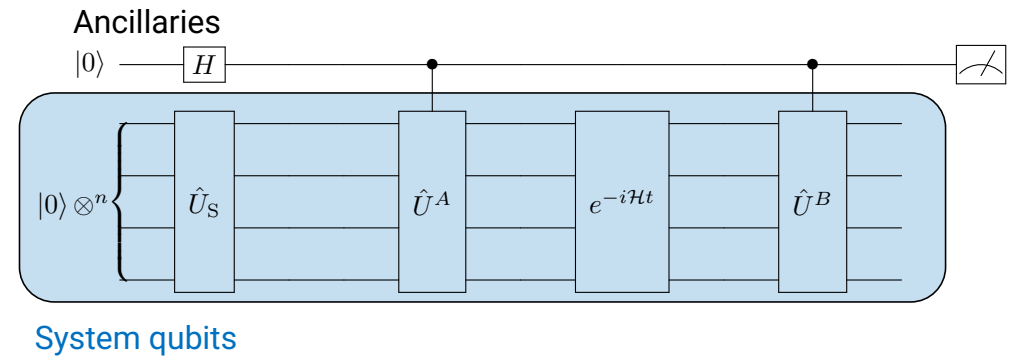
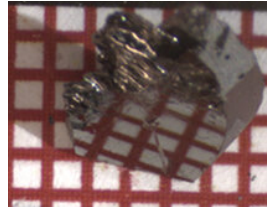
Interfere with ground state Complete expectation value Time evolve Prepare state of interest

Apply excitation B Apply excitation A

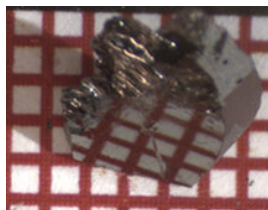
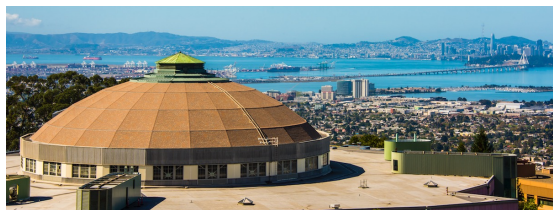


Somma, Simulating physical phenomena by quantum networks (2002)

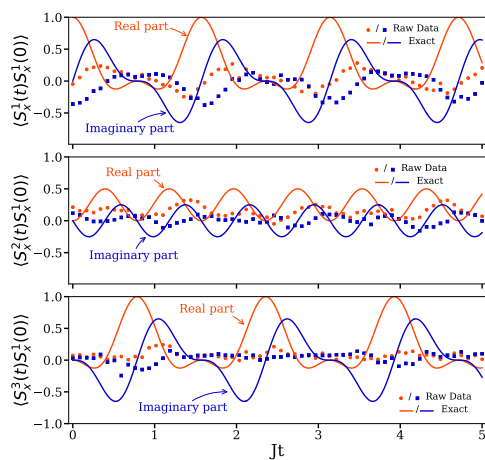
Correlation functions



Correlation functions

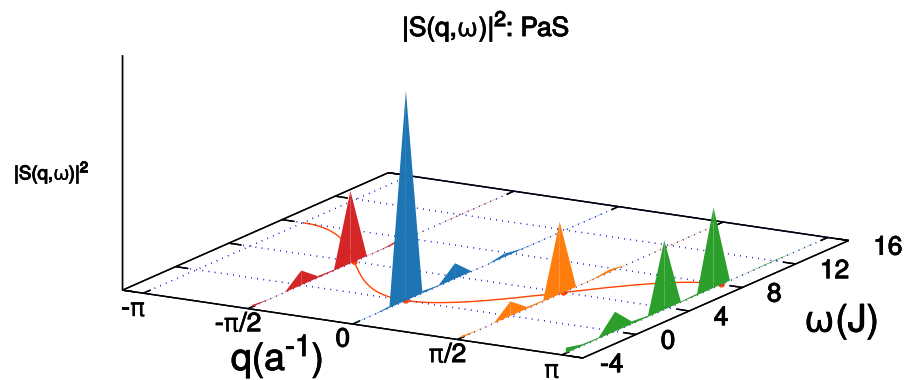
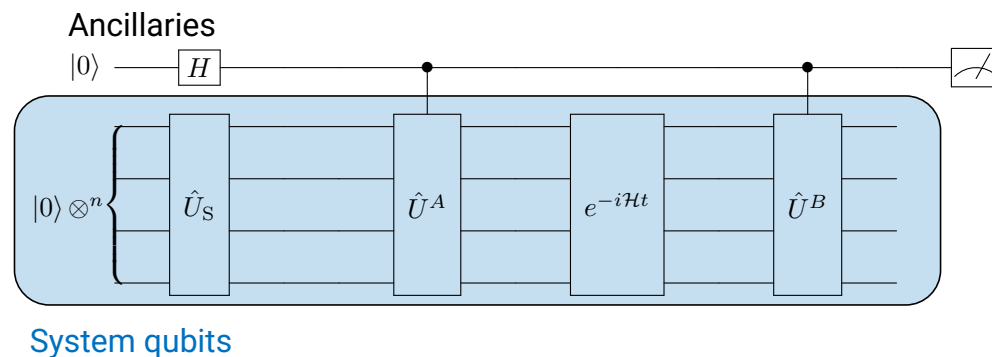
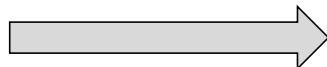


Raw data (2019)

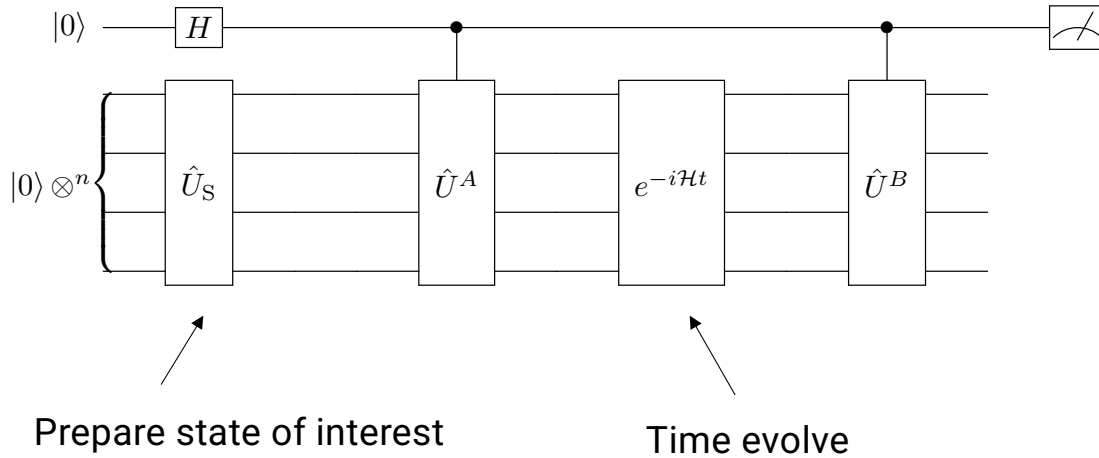


$$\langle A(r, t) B(r', t') \rangle$$

Error mitigation



A-Z quantum simulation



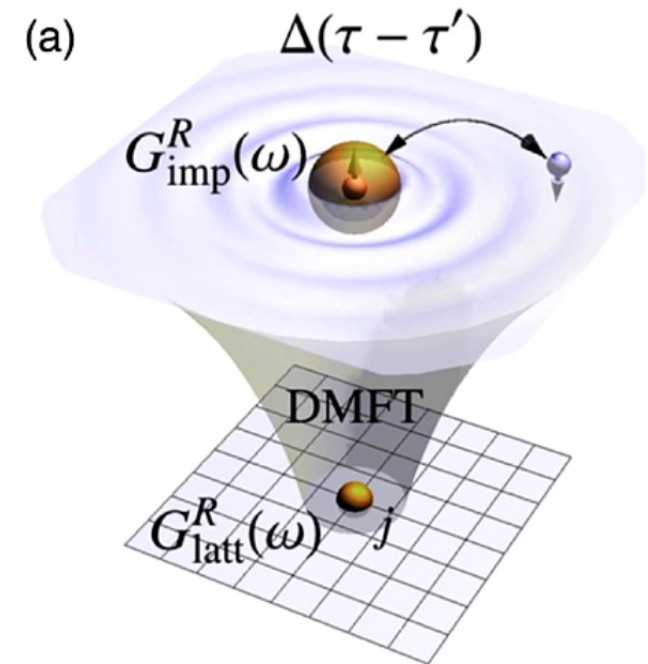
☹️ Circuit to prepare interacting ground state is very deep

☹️ Variational approaches are very difficult in the presence of noise

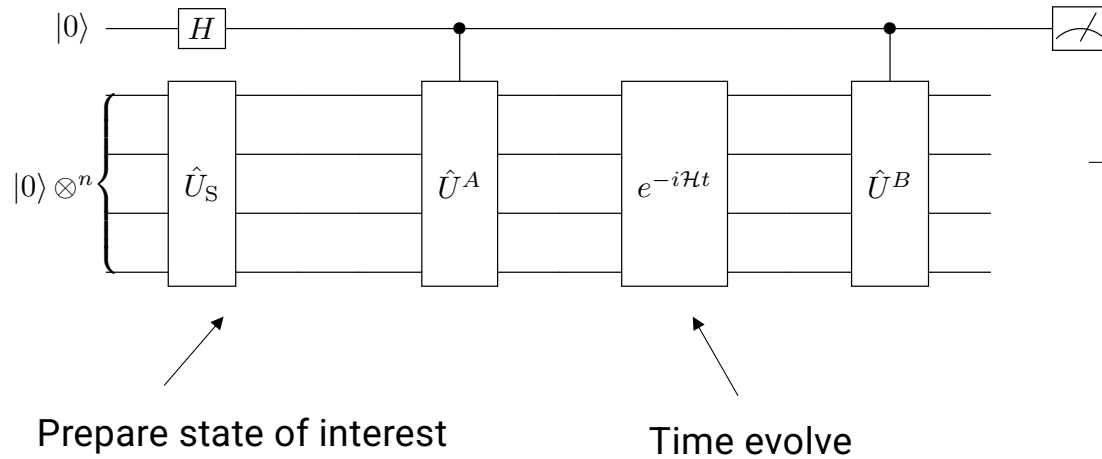
☹️ Standard Trotter decomposition leads to deep circuits with many gates

☹️ Alternative approaches (QSP) requires many ancillae

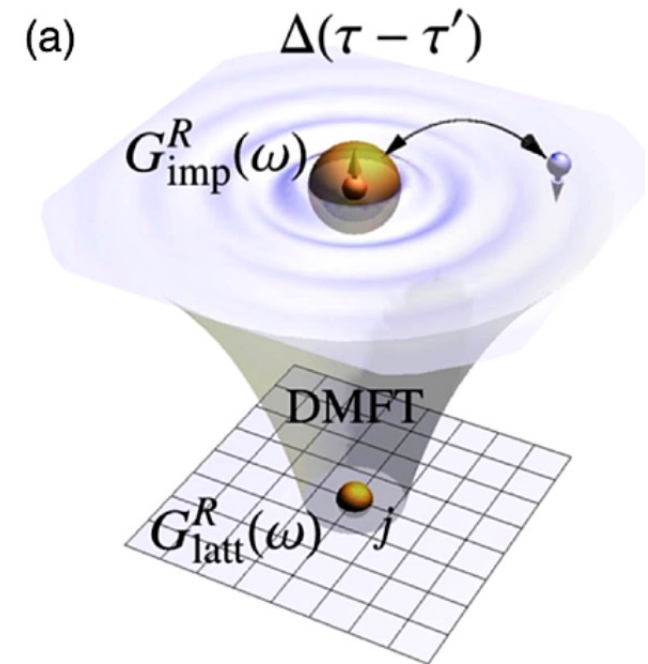
☹️ Your results will be very noisy anyway



A-Z quantum simulation



Classical post-processing techniques



- Physics-Informed Subspace Expansions

- Lie-algebraic methods for time evolution

PHYSICAL REVIEW LETTERS **121**, 032501 (2018)

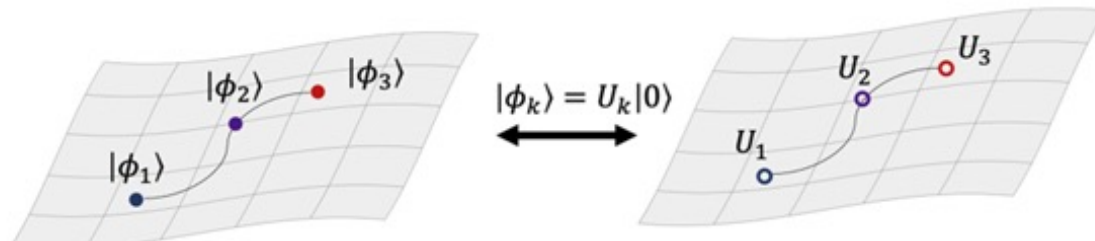
Featured in Physics

Eigenvector Continuation with Subspace Learning

Dillon Frame,^{1,2} Rongzheng He,^{1,2} Ilse Ipsen,³ Daniel Lee,⁴ Dean Lee,^{1,2} and Erma Rrapaj⁵

- Ground state varies continuously in a parameter space and is spanned by a few low energy state vectors.

$$|\phi_3\rangle = \alpha_1 |\phi_1\rangle + \alpha_2 |\phi_2\rangle$$

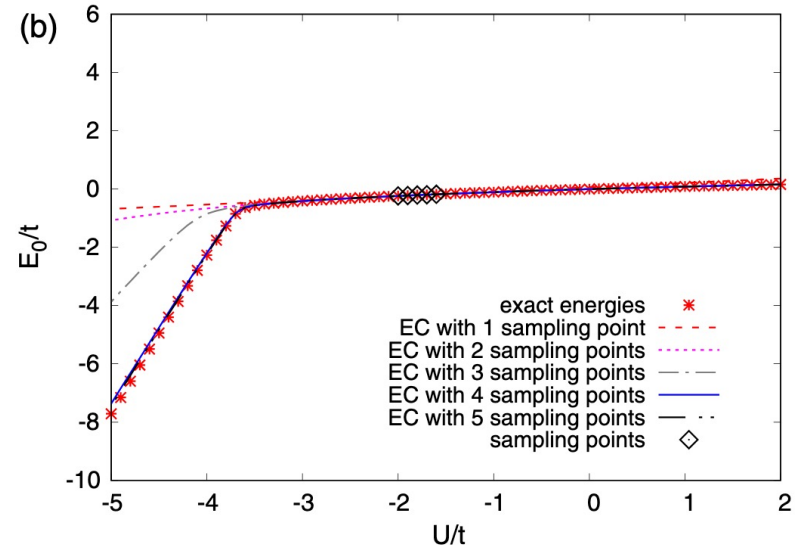
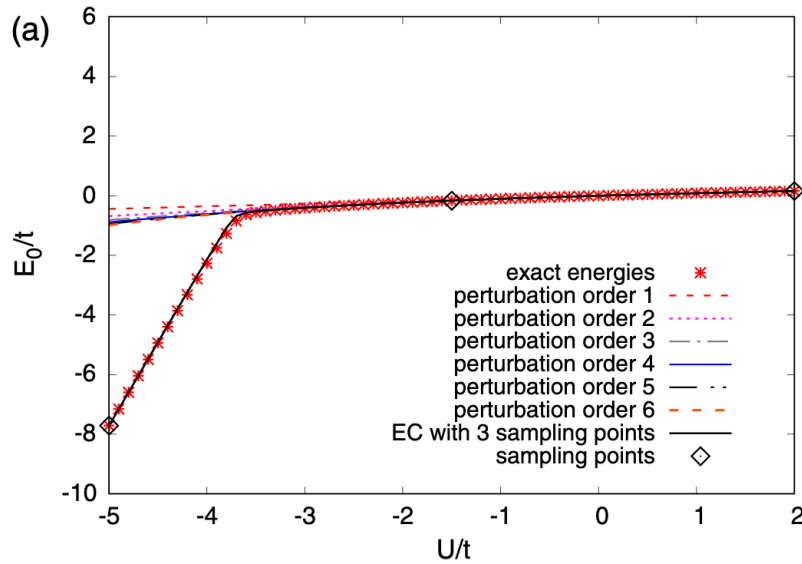


PHYSICAL REVIEW LETTERS **121**, 032501 (2018)

Featured in Physics

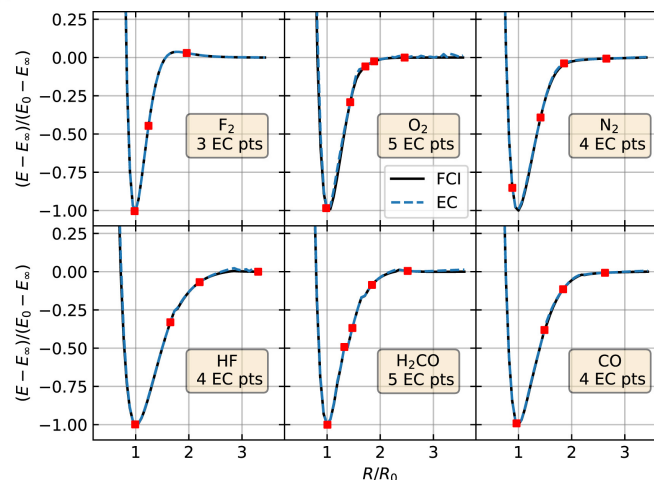
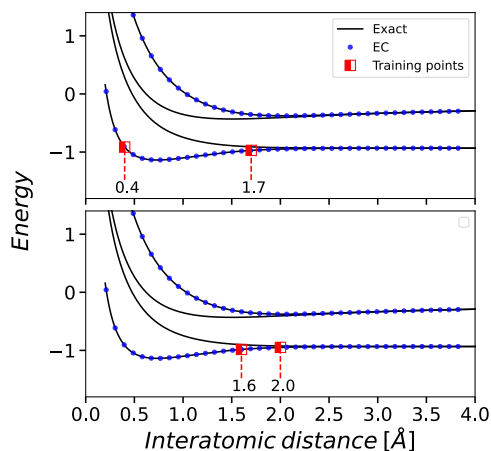
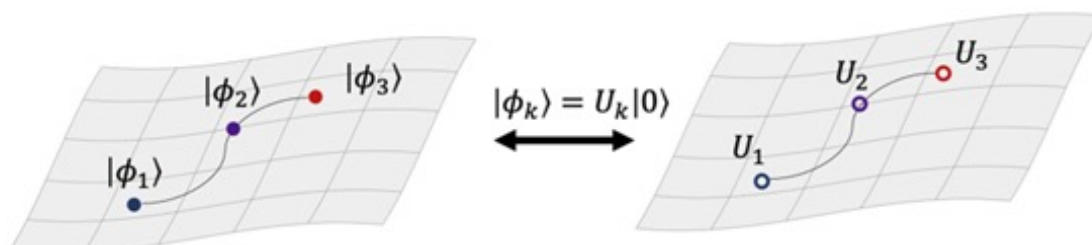
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arXiv:2209.10571

C. Mejuto-Zaera et al., Electron. Struct. 2023

Commun. Math. Phys. 356, 451–500 (2017)
Digital Object Identifier (DOI) 10.1007/s00220-017-2976-9

Communications in
**Mathematical
Physics**



Complexity of Quantum Impurity Problems

Sergey Bravyi, David Gosset

IBM T.J. Watson Research Center, Yorktown Heights, NY, USA. E-mail: sbravyi@us.ibm.com;
dngosset@us.ibm.com

Received: 28 November 2016 / Accepted: 6 June 2017
Published online: 31 August 2017 – © Springer-Verlag GmbH Germany 2017

“The low energy state is represented as a superposition of $\exp [O(b^3)]$ fermionic Gaussian states.”

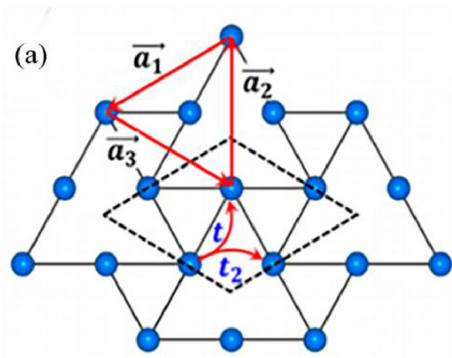
PHYSICAL REVIEW RESEARCH 3, 033188 (2021)

Quantum impurity models using superpositions of fermionic Gaussian states: Practical methods and applications

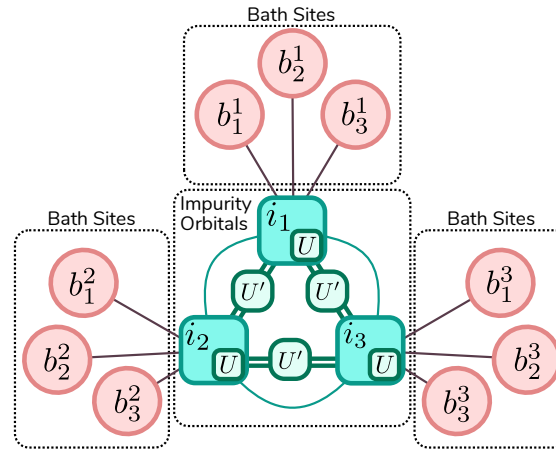
Samuel Boutin and Bela Bauer
Station Q, Microsoft Corporation, Santa Barbara, California 93106 USA

“Quantum impurity models provide a natural arena for studying the complexity of fermionic systems in an intermediate regime interpolating between the free and the fully interacting cases.”

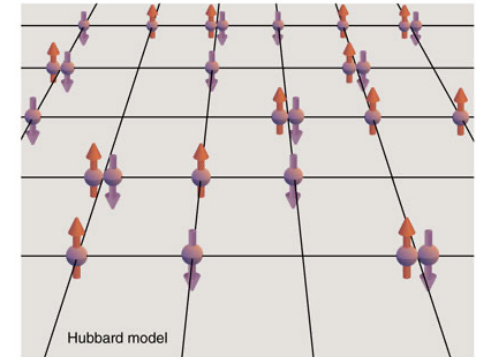
Fermionic Gaussian Subspace



Free fermions



Impurity model



Fully interacting

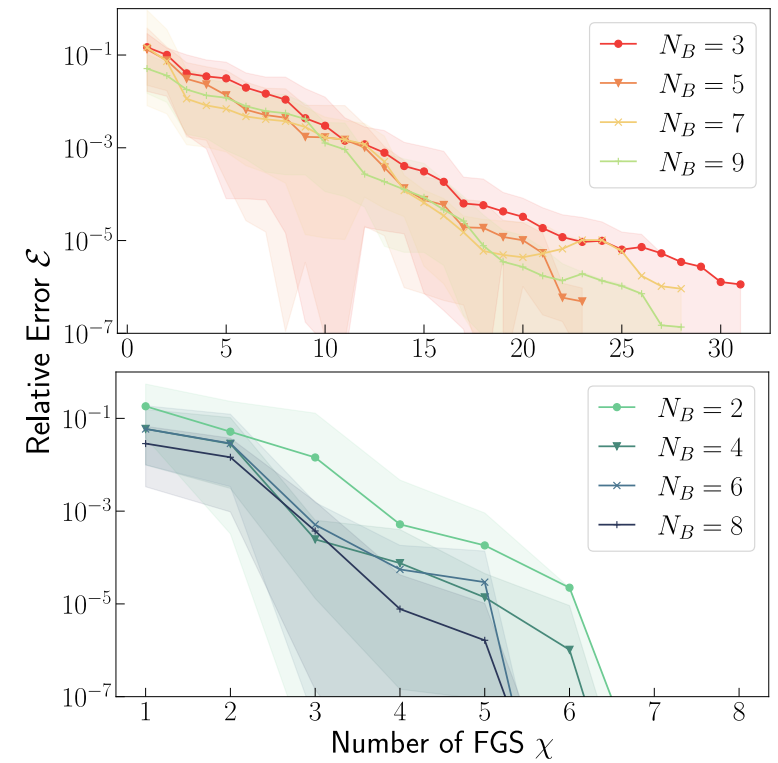
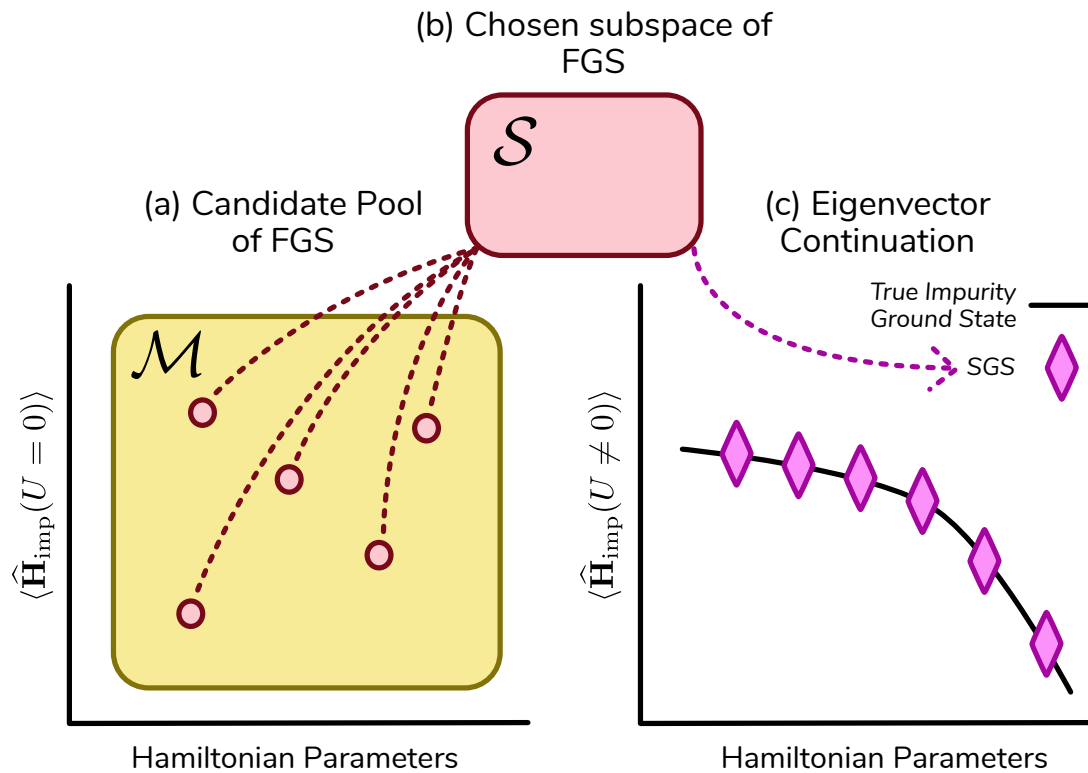
Complexity

Advantages of a basis based on fermionic gaussian states:

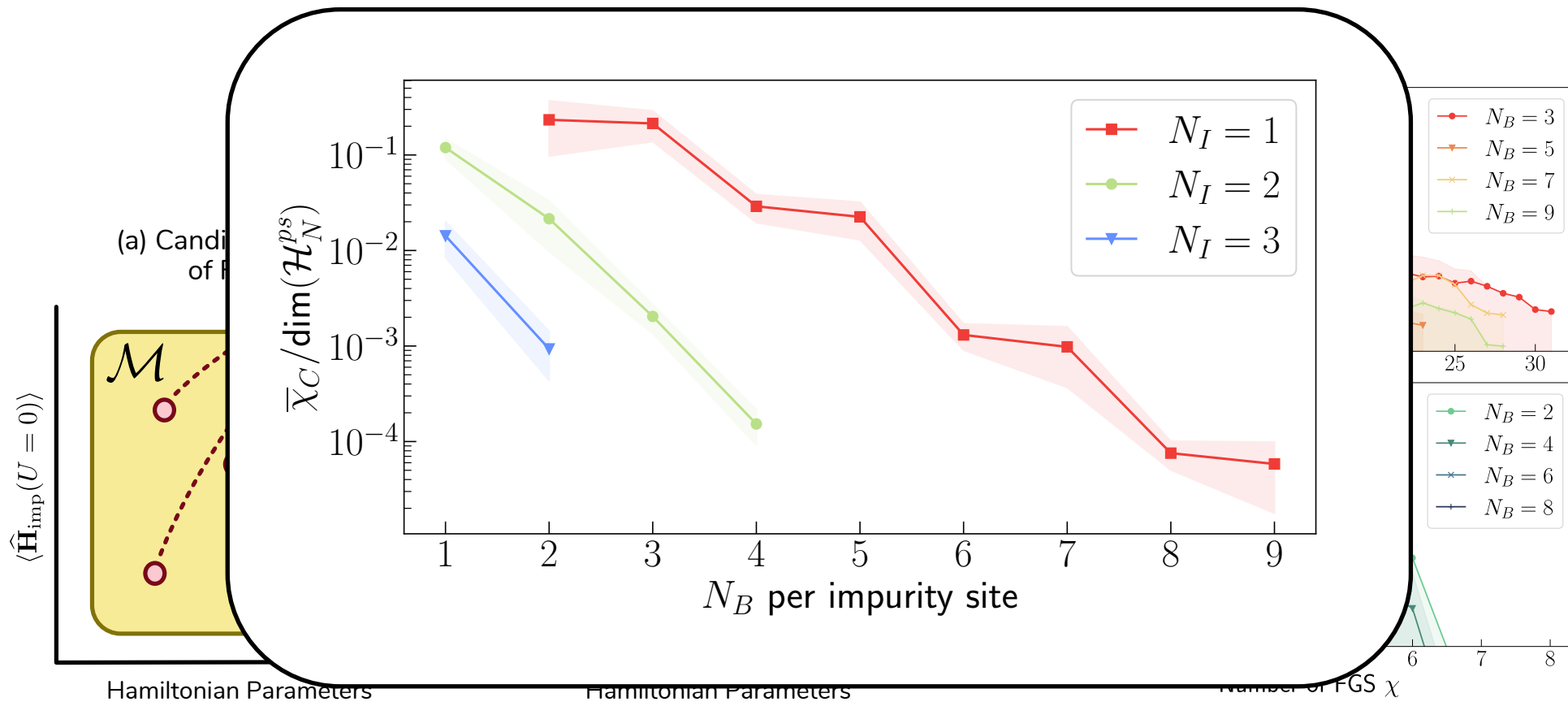
- Polynomially sized calculations to find the ground state
- Easy to prepare on QC
- One basis set spans the necessary space across entire DMFT phase diagram

Fermionic Gaussian Subspace

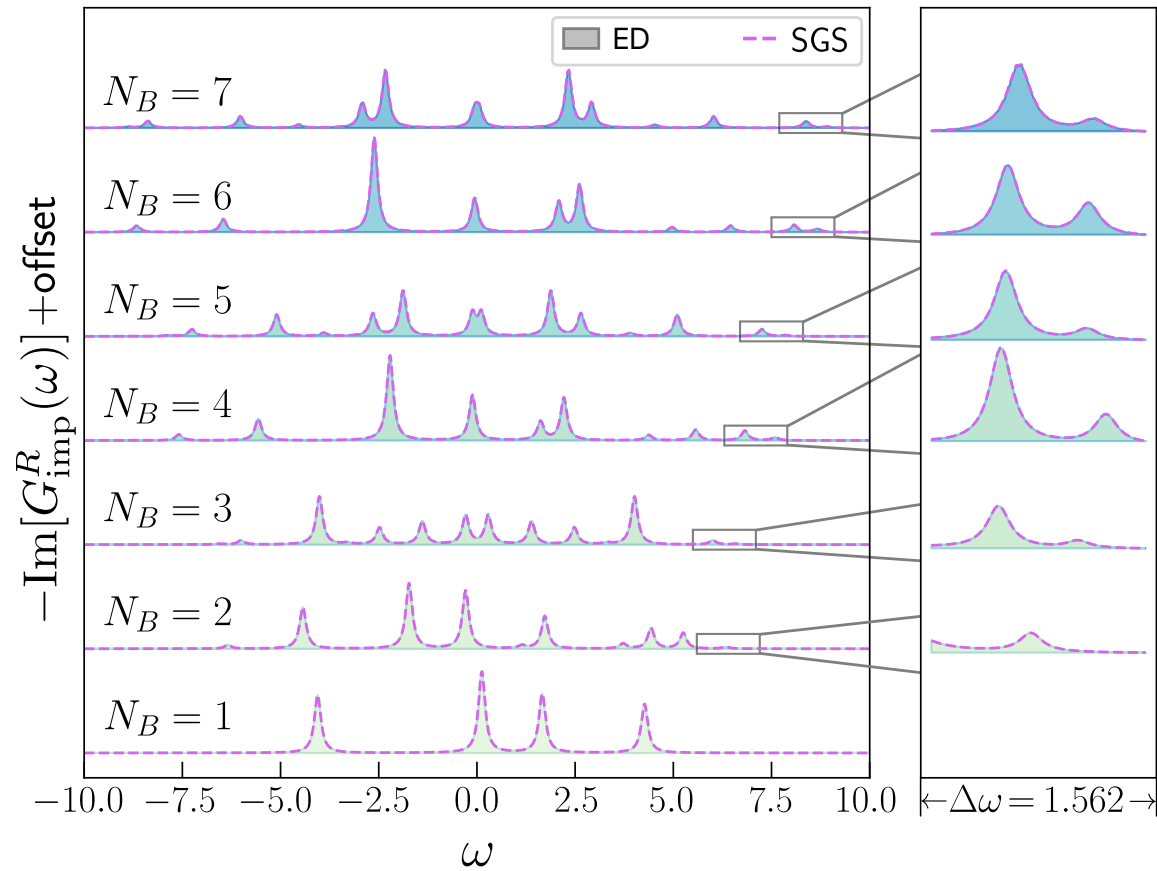
$$\tilde{H}\vec{\alpha} = \tilde{E}S\vec{\alpha}$$



Fermionic Gaussian Subspace



Fermionic Gaussian Subspace



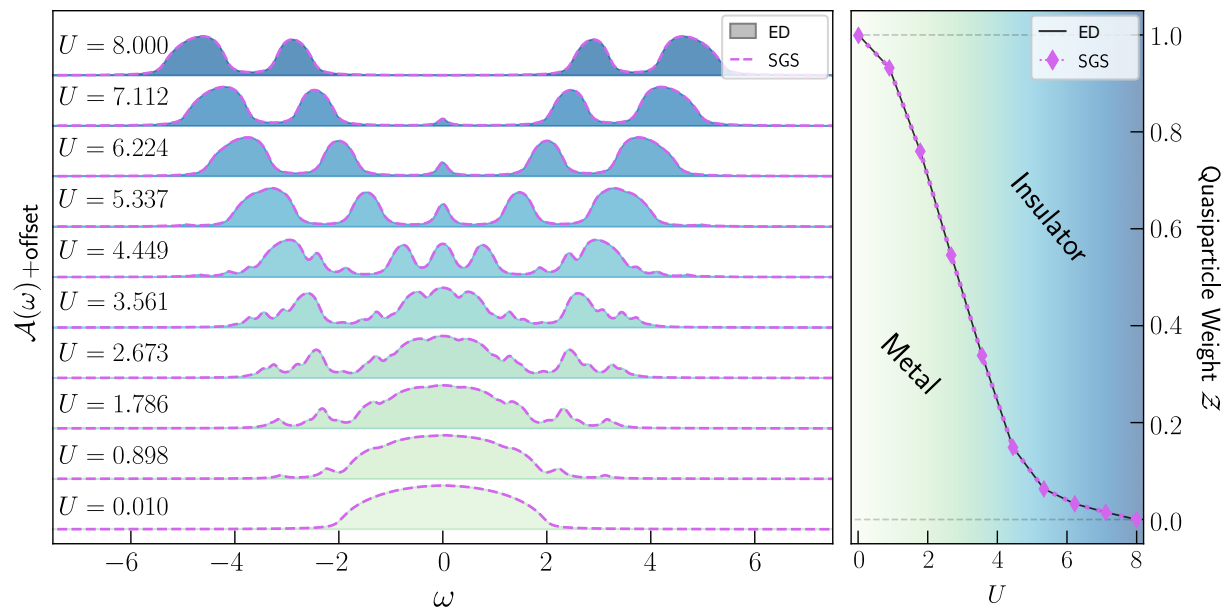
Fermionic Gaussian Subspace

Now that we have the tools, let's see how it works:

- DMFT using sum of Gaussian states (SGS) (Hubbard Model w/ Bethe lattice)

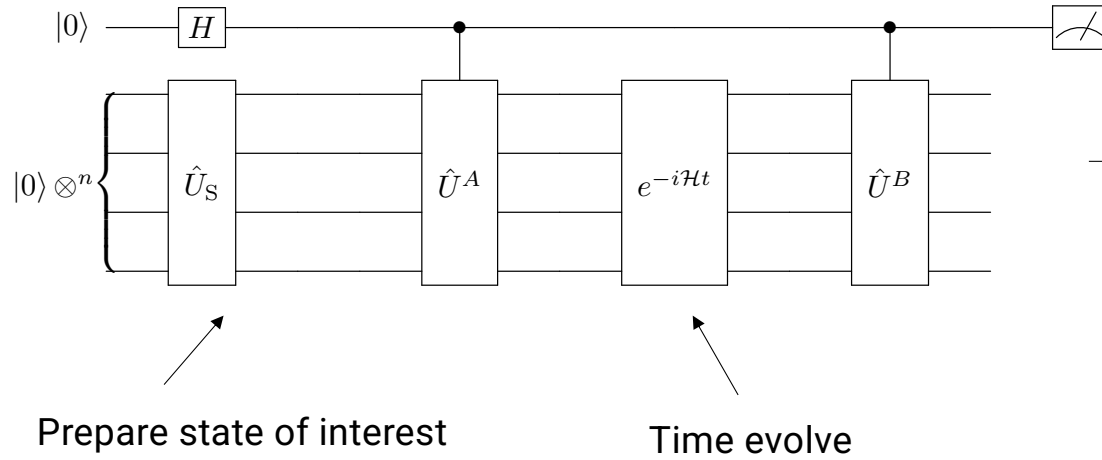
Shaded region = $|\psi\rangle$

$$\text{dashed line} = \sum_{k=1}^{\chi} \alpha_k |\phi_k\rangle$$



Increasing
interaction
strength on
impurity

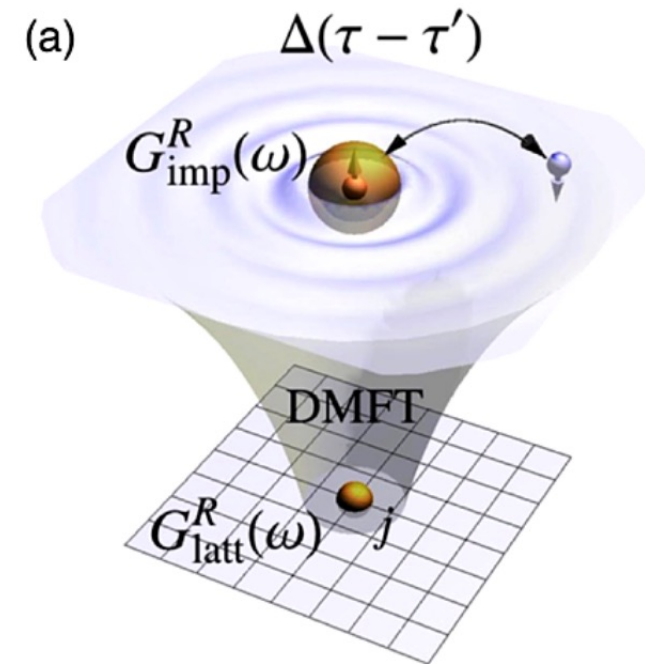
A-Z quantum simulation



Classical post-processing techniques

- *Physics-Informed Subspace Expansions*

- *Lie-algebraic methods for time evolution*

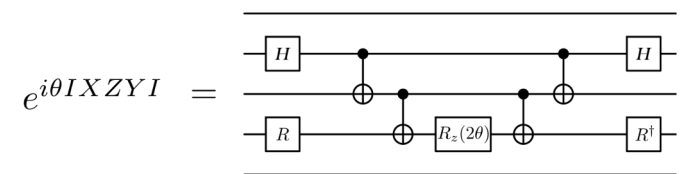
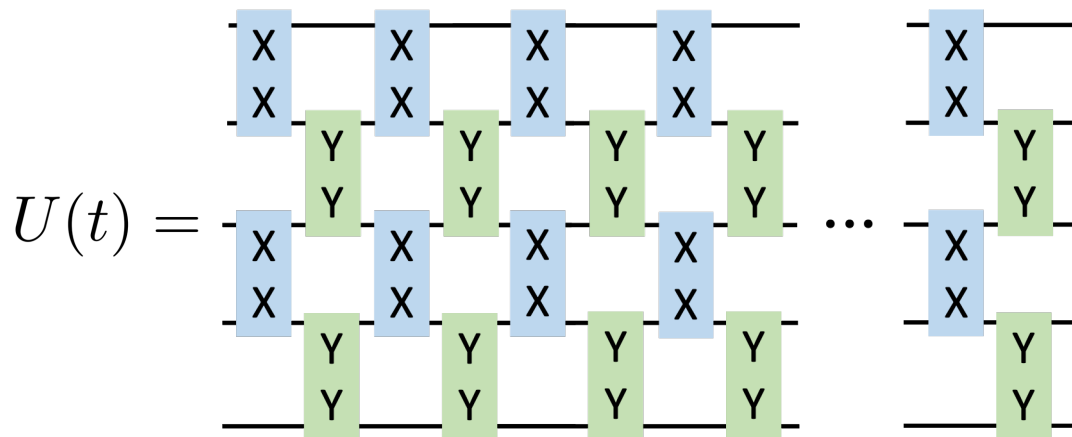


Time evolution

Simulation of a time independent spin Hamiltonian:

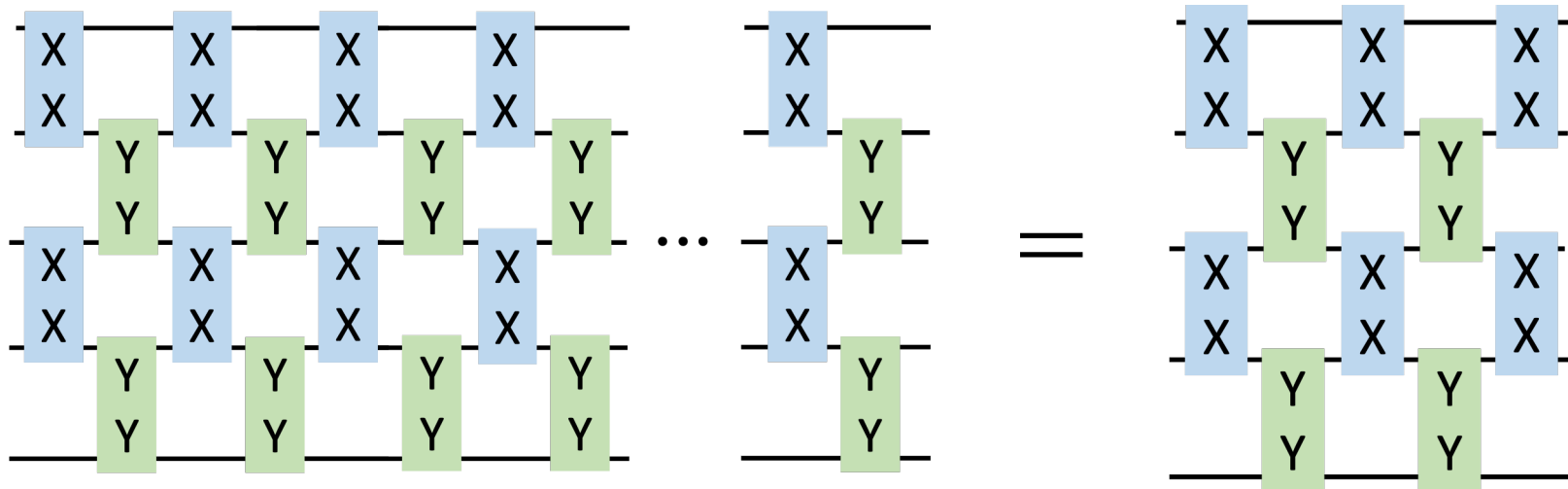
$$\mathcal{H} = a XXIII + b IYYII + c IIXXI + d IIIYY$$

$$U(\epsilon) = e^{-i\epsilon\mathcal{H}} = e^{-i\epsilon a XXIII} e^{-i\epsilon b IYYII} e^{-i\epsilon c IIXXI} e^{-i\epsilon d IIIYY} + O(\epsilon^2)$$



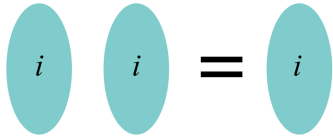
Algebraic Circuit Compression

- We propose a **constructive**, Lie algebra based method which leads to fixed depth circuits for several models
- The method is **scalable** due to its “constructive” and “local” nature.

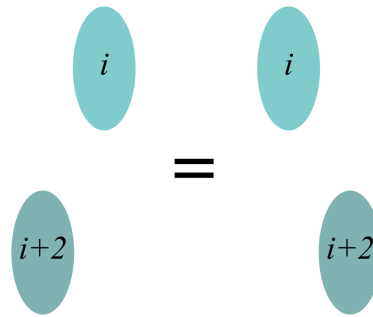


We define an abstract object called “block” which satisfies:

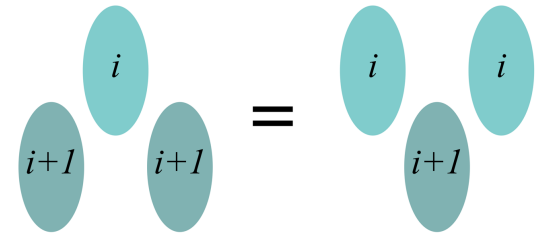
Fusion



Commutation

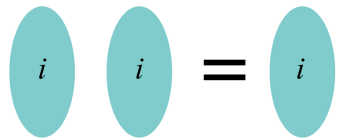


Turnover

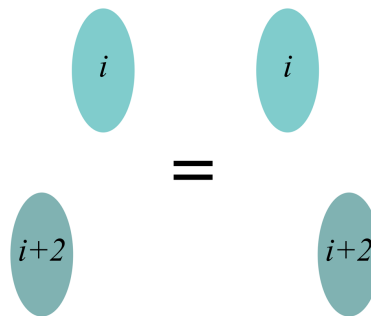


Algebraic Circuit Compression

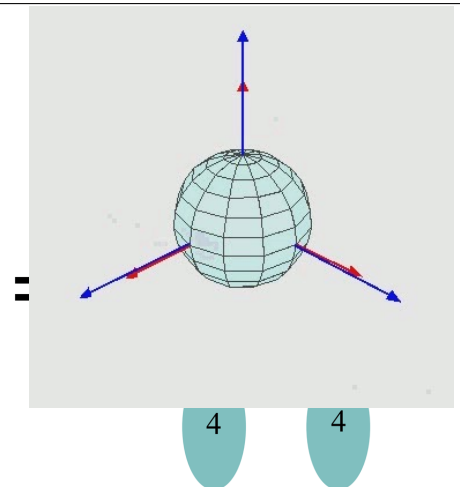
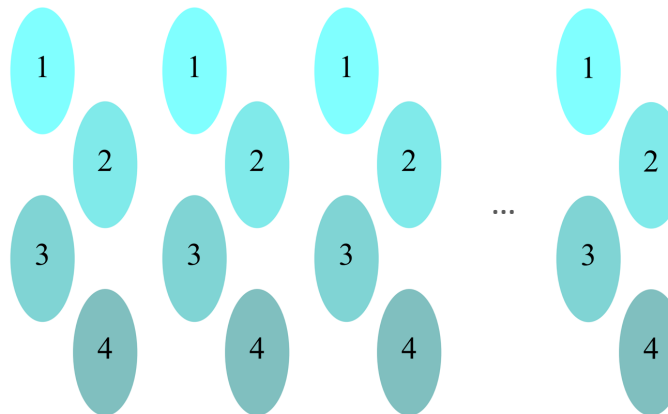
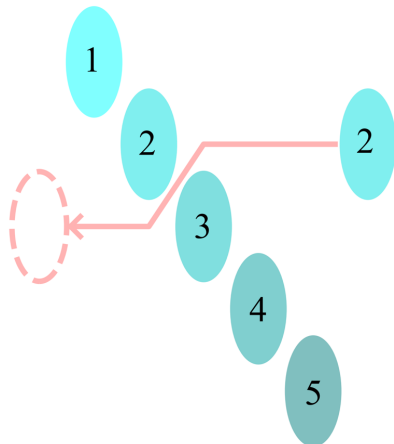
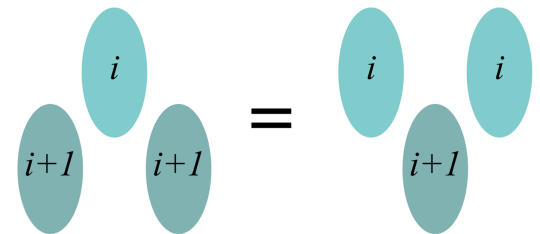
Fusion



Commutation



Turnover



Free Fermion Blocks

Kitaev Chain

$$2i-1 = \begin{array}{c} 2i-1 \\ 2i \end{array} \begin{array}{c} \text{X} \\ \text{X} \end{array}$$

$$2i = \begin{array}{c} 2i \\ 2i+1 \end{array} \begin{array}{c} \text{Y} \\ \text{Y} \end{array}$$

Transverse Field Ising

$$2i-1 = i \text{---} \text{Z} \text{---}$$

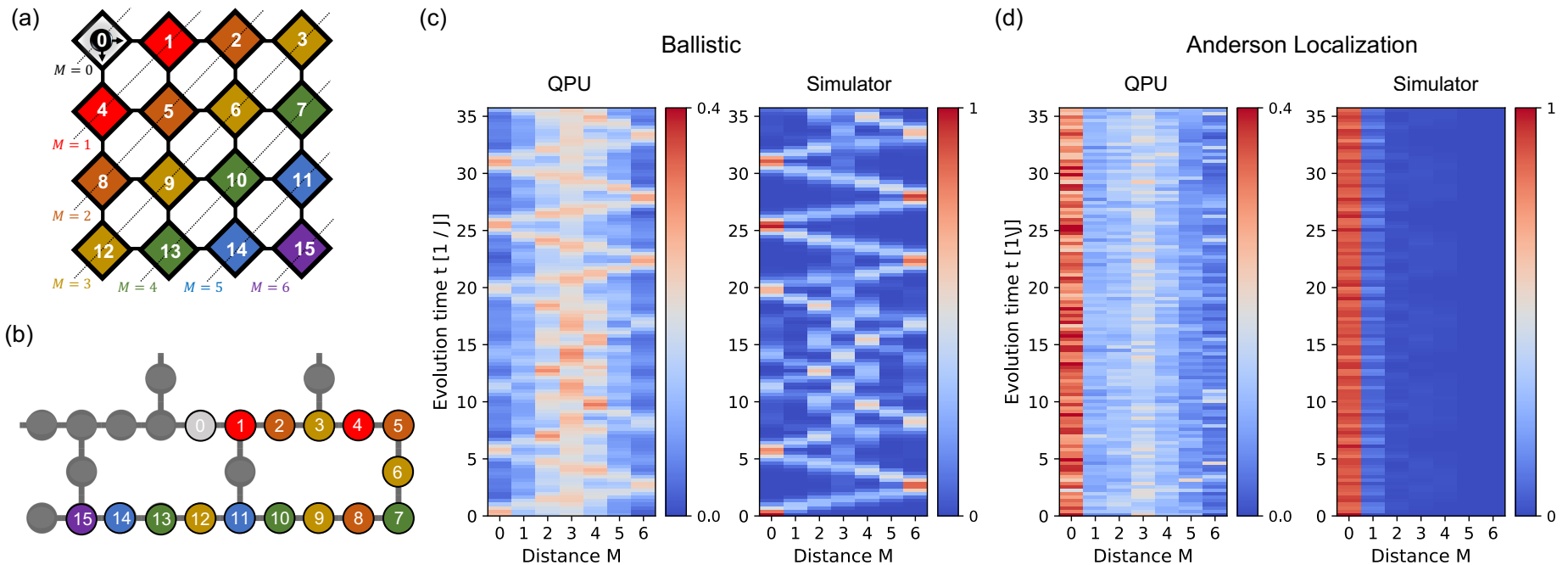
$$2i = \begin{array}{c} i \\ i+1 \end{array} \begin{array}{c} \text{X} \\ \text{X} \end{array}$$

Transverse Field XY

$$i = \begin{array}{c} i \\ i+1 \end{array} \begin{array}{c} \text{Z} \\ \text{Z} \end{array} \begin{array}{c} \text{X} \\ \text{X} \end{array} \begin{array}{c} \text{Y} \\ \text{Y} \end{array} \begin{array}{c} \text{Z} \\ \text{Z} \end{array}$$

$$\mathcal{H} = -t \sum_i [X_i X_{i+1} + Y_i Y_{i+1}] + h \sum_i Z_i$$

Approach #1: Algebraic Circuit Compression (preliminary)



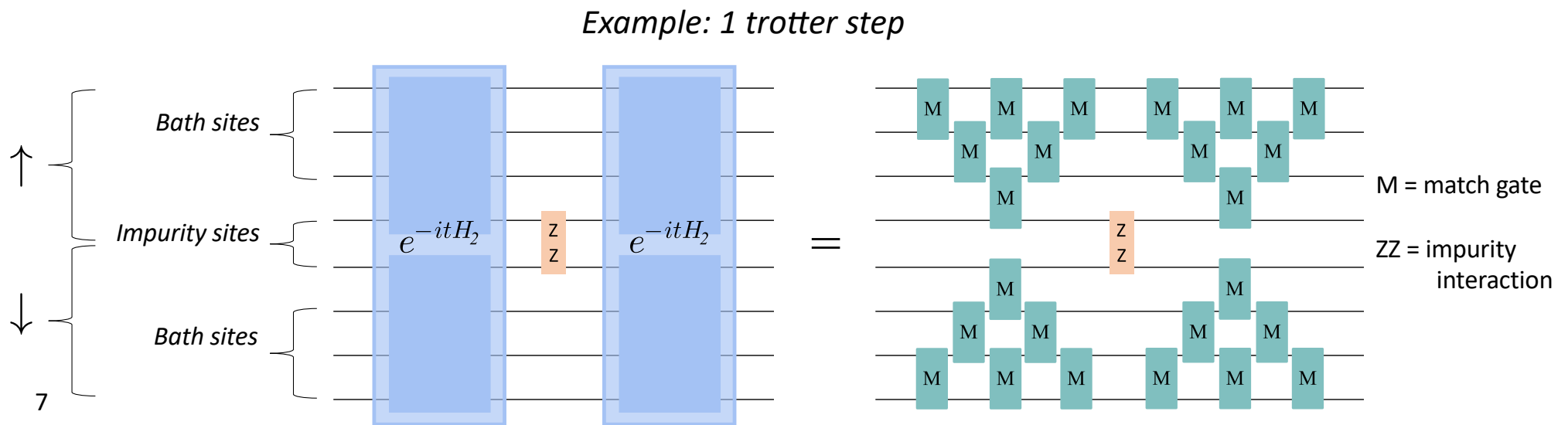
Simulating the Impurity Model on a Quantum Computer

Time evolving the impurity ground state

[1] 10.1103/PhysRevLett.129.070501

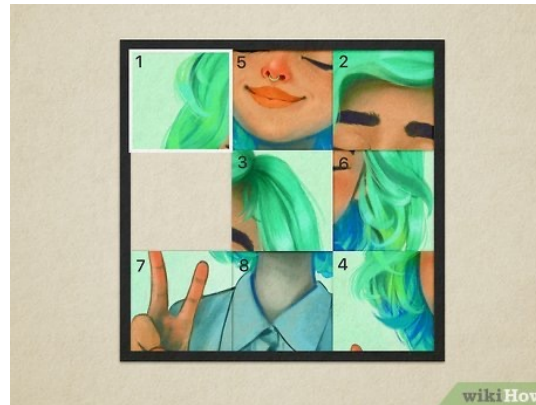
[2] arXiv:2303.09538

➤ **Partial compression** for impurity models:



- [1] 10.1103/PhysRevLett.129.070501

- [2] arXiv:2303.09538

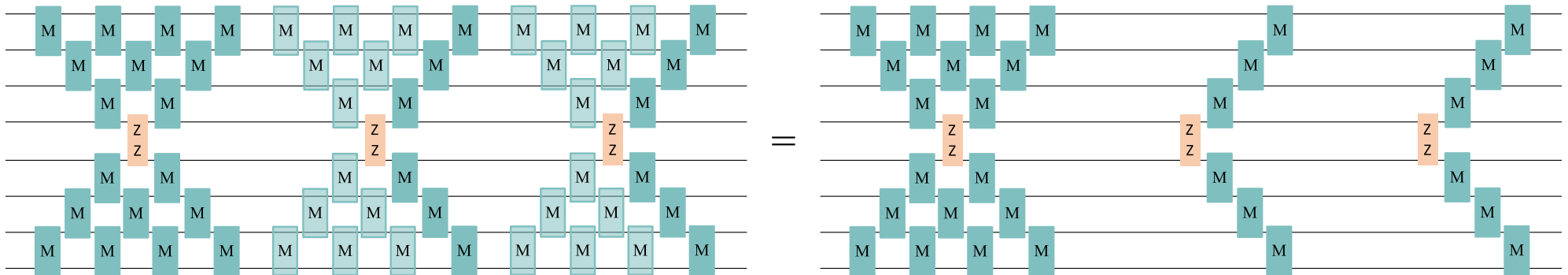


Simulating the Impurity Model on a Quantum Computer

[1] 10.1103/PhysRevLett.129.070501

[2] arXiv:2303.09538

Further example: 3 trotter steps

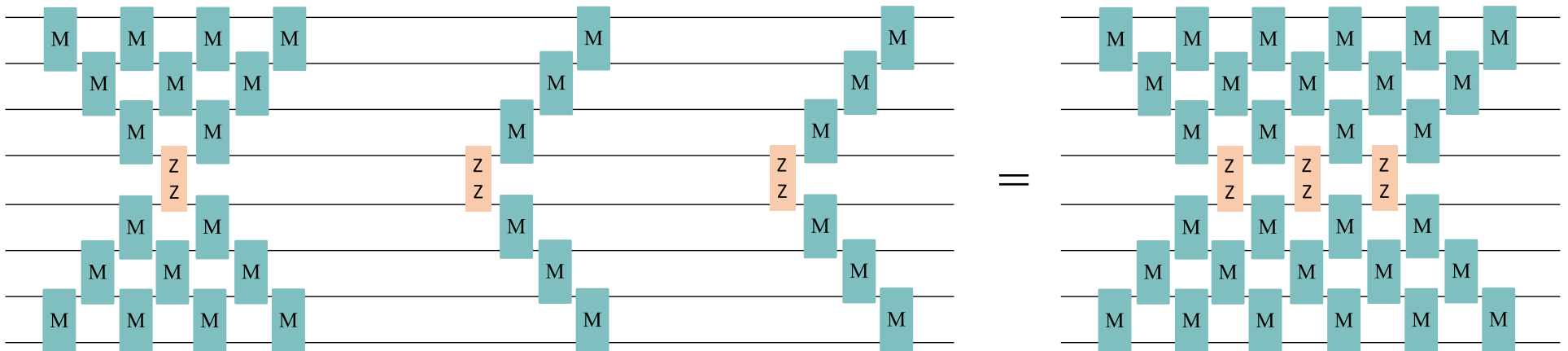


Simulating the Impurity Model on a Quantum Computer

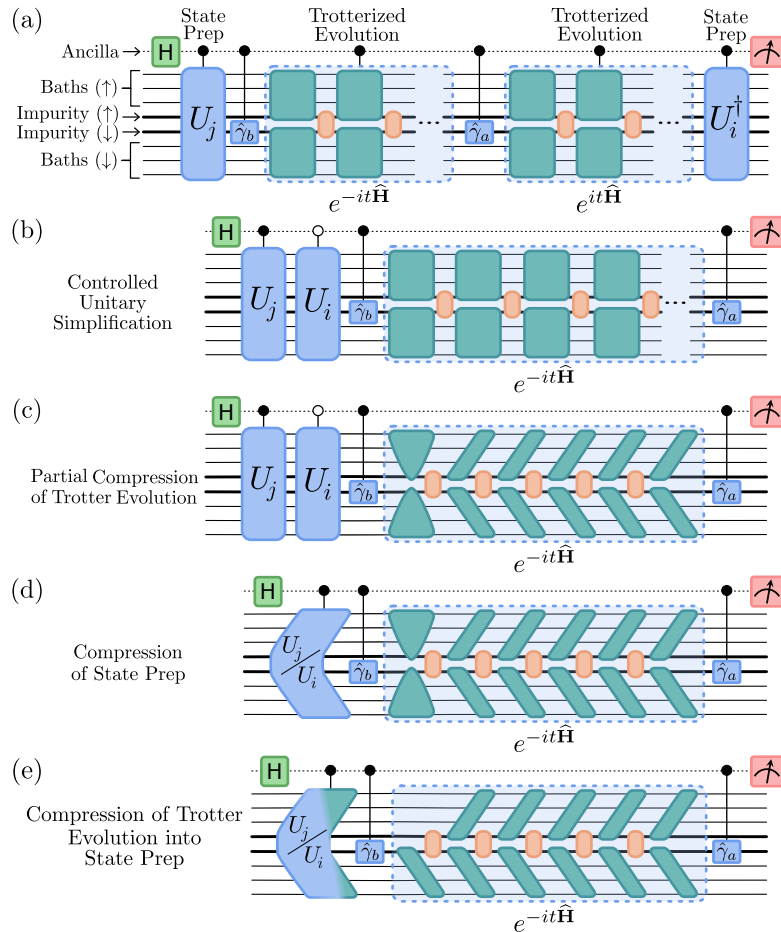
[1] 10.1103/PhysRevLett.129.070501

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Further example: 3 trotter steps

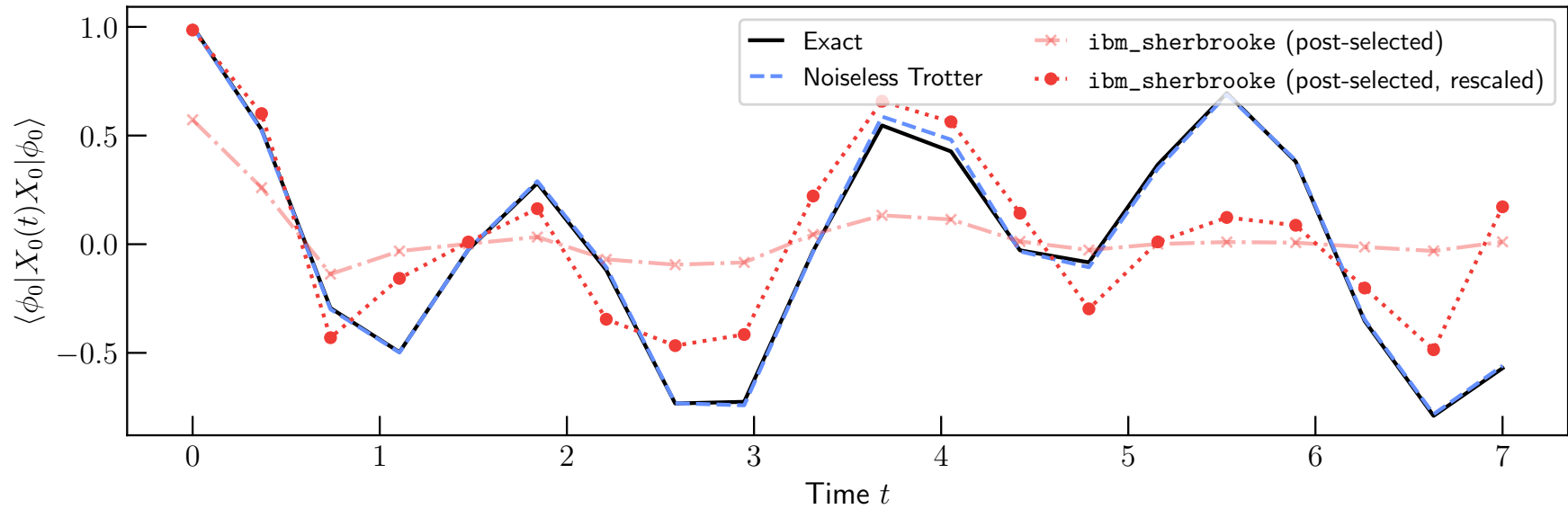


Simulating the Impurity Model on a Quantum Computer



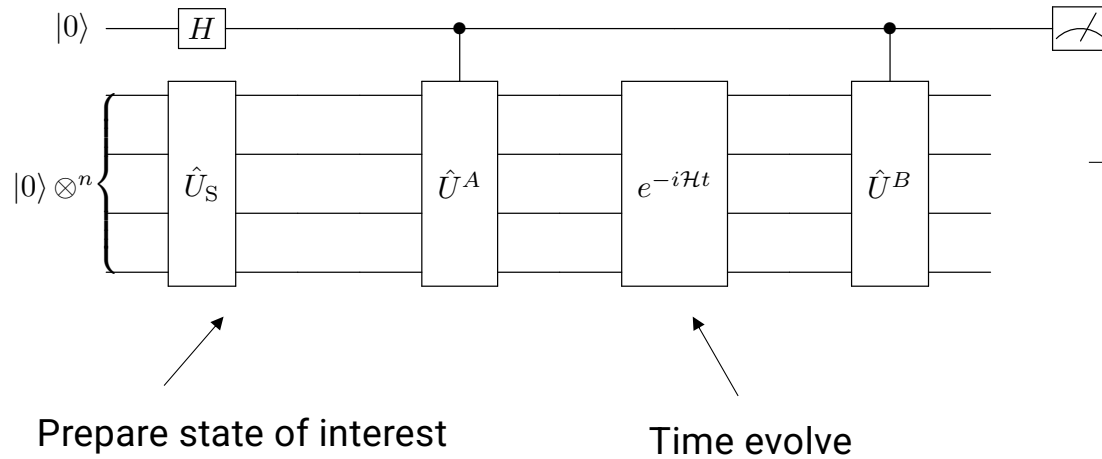
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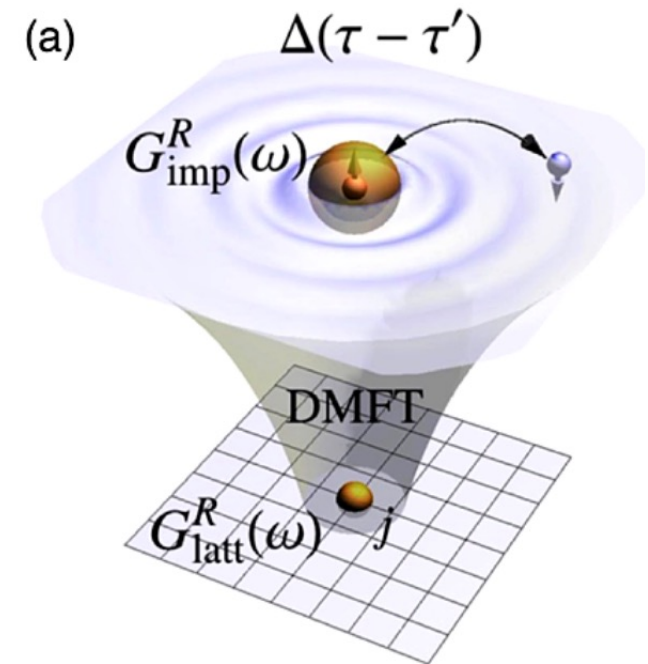
$$\mathcal{G}_{00}(t) = -i\theta(t)\langle\phi_0|\{c_0(t), c_0^\dagger\}|\phi_0\rangle$$

A-Z quantum simulation



- *Physics-Informed Subspace Expansions*

- *Lie-algebraic methods for time evolution*



A-Z quantum simulation

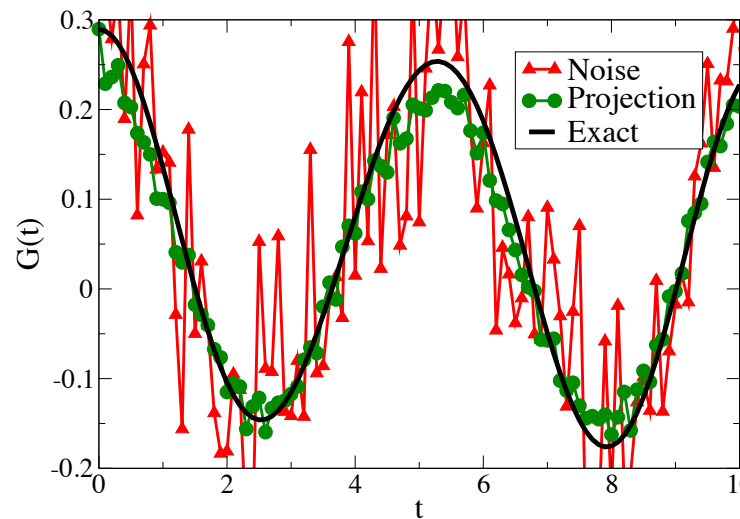
- It turns out that these are positive semi-definite (PSD) functions:

$$G_{AA}(t - t') = \text{Tr} [\rho A(t)^\dagger A(t')]$$

- Then this is a PSD matrix:

$$\underline{G} = \begin{pmatrix} f_0 & f_1 & f_2 & \cdots & f_n \\ f_1^* & f_0 & f_1 & \cdots & f_{n-1} \\ f_2^* & f_1^* & f_0 & \cdots & f_{n-2} \\ \vdots & & & \ddots & \vdots \\ f_n^* & f_{n-1}^* & f_{n-2}^* & \cdots & f_0 \end{pmatrix}$$

where $G_{AA}(t_i - t_j) \rightarrow f_{i-j}$



A-Z quantum simulation

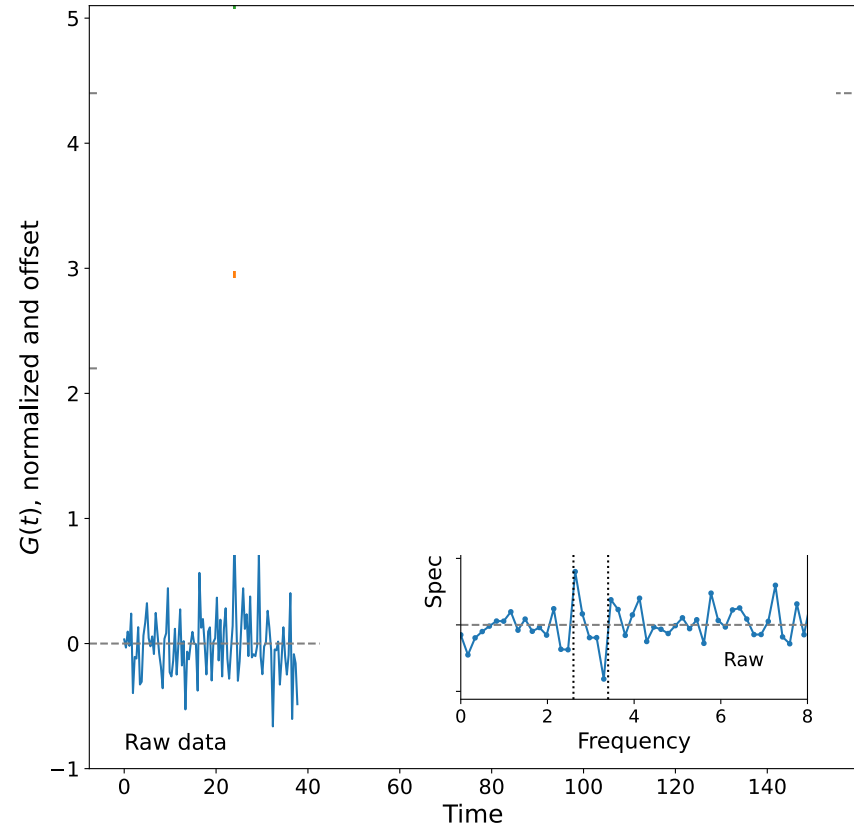
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A-Z quantum simulation

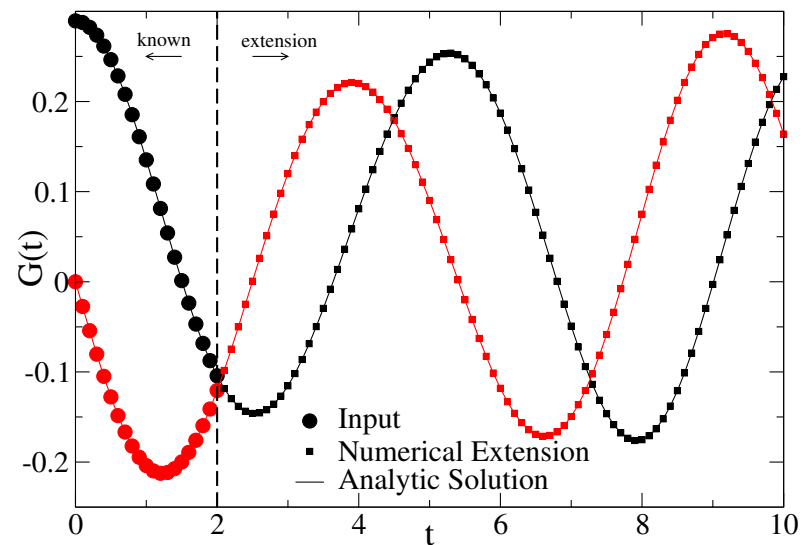
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A-Z quantum simulation

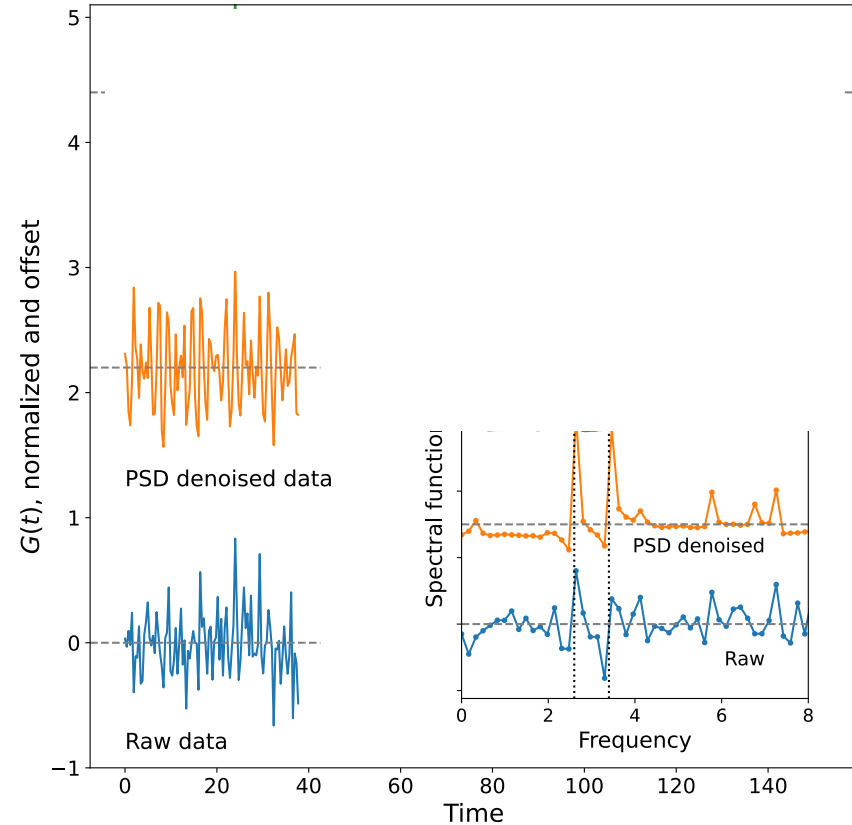
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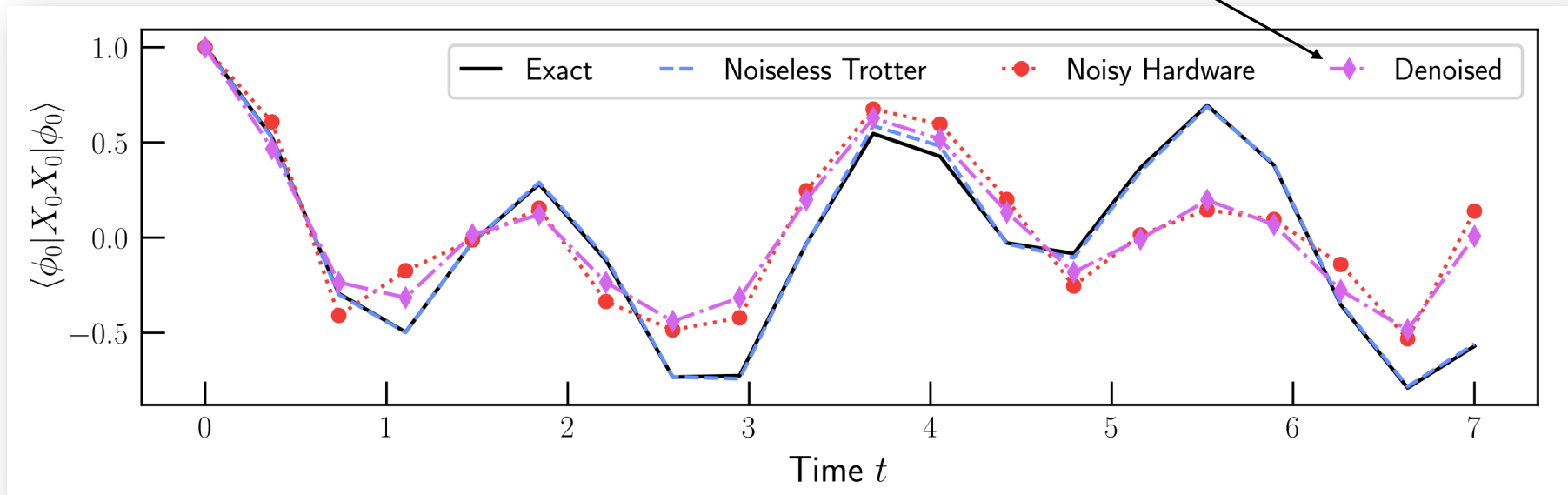
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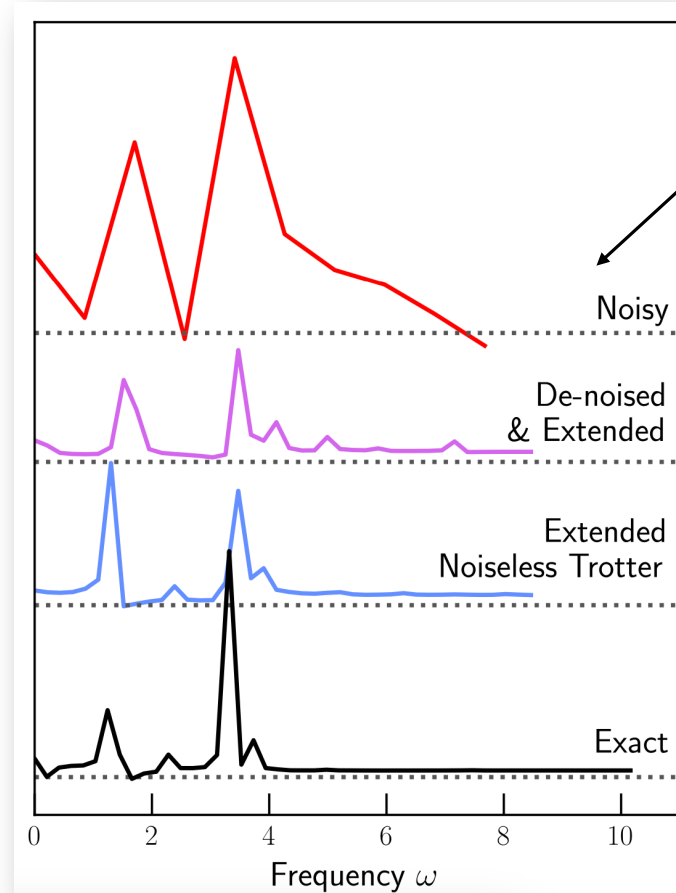
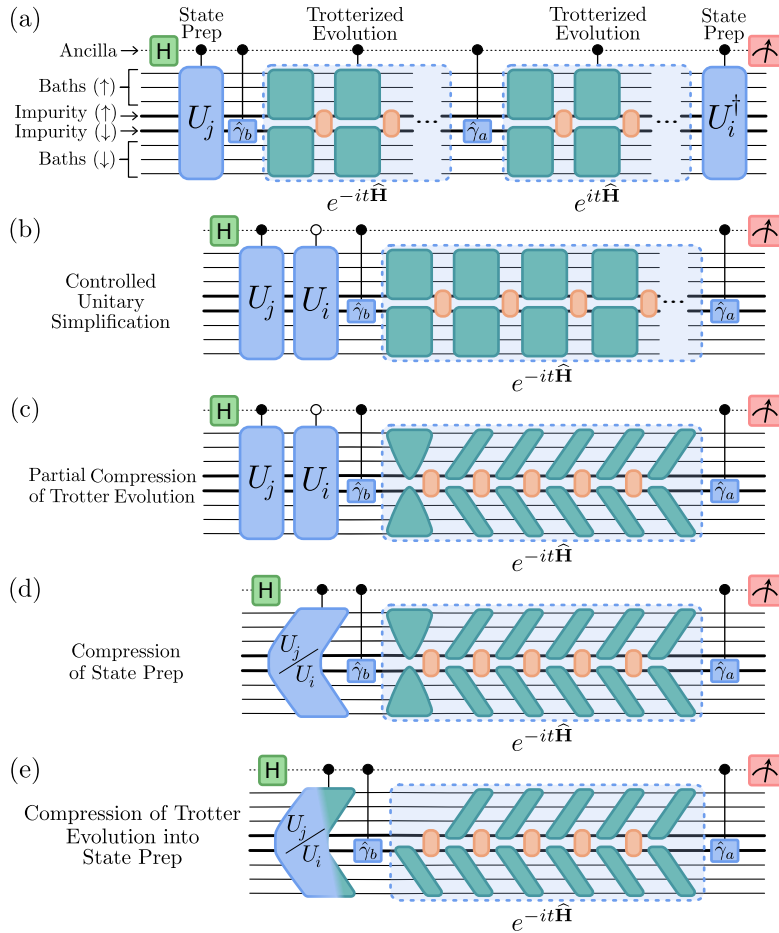


Hardware results

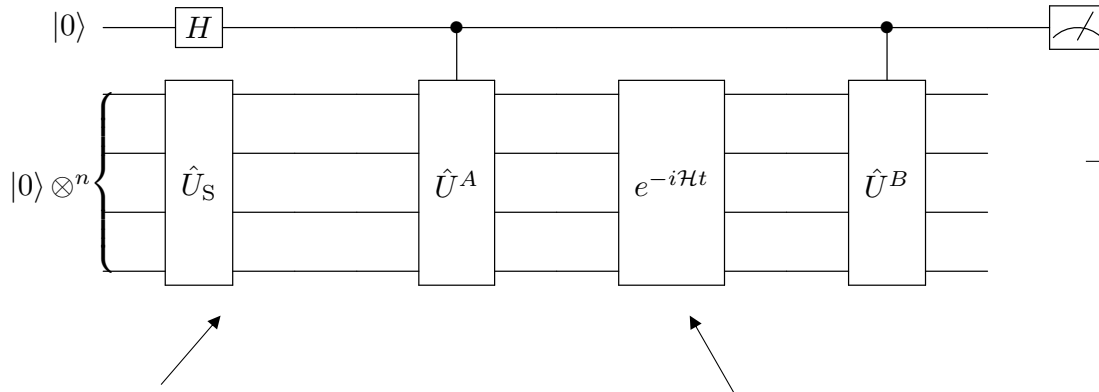
- Evaluation of correlation functions on a quantum computer (ibm_sherbrooke)
 - Shallow circuits + error mitigation = signal we can work with
 - Signal-processing used (post-selecting results, PSD de-noising [10.1103/PhysRevLett.132.160403])



Hardware results



DMFT on Quantum Computers – a path to quantum advantage



Classical post-processing
technique based on PSD
projection

Prepare state of interest
using subspace of free
fermionic states

Time evolve
using circuit compression

